This is a mathematical intermezzo on the "undetermined multiplier method of Lagrange".

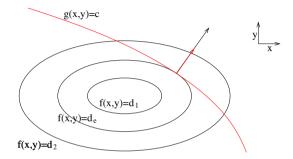
The aim: minimizing a function (e.g. the expectation value of H) while applying a constraint (e.g. the orbitals in the wave function should be orthogonal).

First a graphical explanation for a simple case.

Wanted: extremum of f(x, y) subject to the constraint: g(x, y) = c

g(x, y) = c defines a curve, i.e. a single contour of g(x, y) with g(x, y) = c.

We look at contours of f(x, y), i.e. curves for which f(x, y) = d, with d a constant. Near the extremum there are 2 or no intersections of g(x, y) = c and f(x, y) = d, depending on the value of d.



We change d, from d_1 to d_2 (see plot), and, when we hit the extremal point (x_e, y_e) , the two intersections merge into a single point. In this point, the "surfaces" defined by g(x, y) = c and $f(x, y) = d_e$ are parallel, hence their normal vectors, i.e. the gradients, are parallel:

$$\nabla f = \lambda \nabla g$$
 i.e. $\begin{pmatrix} \partial f / \partial x \\ \partial f / \partial y \end{pmatrix} = \lambda \begin{pmatrix} \partial g / \partial x \\ \partial g / \partial y \end{pmatrix}$

Intermezzo in Intermezzo: The Gradient.

Let's choose a parametrized trajectory s(t) on the constraint:

 $g(s_{\!\scriptscriptstyle X}(t),s_{\!\scriptscriptstyle Y}(t))=c$

Now differentiate:

$$\frac{d}{dt}g(s_x(t),s_y(t)) = \frac{\partial g}{\partial x}\frac{ds_x}{dt} + \frac{\partial g}{\partial y}\frac{ds_y}{dt} = \begin{pmatrix} \frac{\partial g}{\partial x} \\ \frac{\partial g}{\partial y} \end{pmatrix} \cdot \begin{pmatrix} \frac{ds_x}{dt} \\ \frac{ds_y}{dt} \end{pmatrix} = 0$$

The first vector is the gradient.

The second vector is tangent to the surface g(x, y) = c (Dutch: devector raakt aan het oppervlak).

Hence the gradient of g is perpendicular to the isosurface of g. This is a general statement.

End of Intermezzo in Intermezzo.

Solving the system:

$$\begin{cases} \nabla f = \lambda \nabla g \\ g(x, y) = c \end{cases}$$

yields x_e , y_e and λ . These equations can be obtained by demanding that the gradient of the function $\mathcal{L}(x, y, \lambda)$,

$$\mathcal{L} = f(x, y) - \lambda (g(x, y) - c)$$

vanishes in the three dimensional x, y, λ space:

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abla} g = 0$$

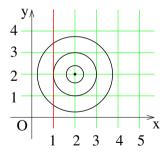
 $\partial \mathcal{L} / \partial \lambda = g(x, y) - c = 0$

Note that we are only guaranteed to get a stationary/critical point, but not necessarily a (local) minimum or maximum.

Let's look at an example ...

$$\begin{cases} f = (x-2)^2 + (y-2)^2 \\ g = x = 1 \end{cases} \ \ \, \text{so} \ \ \, \mathcal{L} = (x-2)^2 + (y-2)^2 - \lambda(x-1)$$

We immediately see that (x, y) = (1, 2) yields the minimum f = 1 on g.



Lagrange:

$$\frac{\partial \mathcal{L}}{\partial x} = 2(x-2) - \lambda = 0$$
$$\frac{\partial \mathcal{L}}{\partial y} = 2(y-2) = 0 \Rightarrow y = 2$$
$$\frac{\partial \mathcal{L}}{\partial \lambda} = -(x-1) = 0 \Rightarrow x = 1$$

 $(x, y, \lambda) = (1, 2, -2)$ is a stationary point in 3D space. Let's look at the "magic" in more detail . . .

Let's try to systematically "complete the square":

$$\mathcal{L} = (x-2)^2 + (y-2)^2 - \lambda(x-1) = x^2 - 4x + 4 - \lambda x + \lambda + (y-2)^2 = \left(x-2-\frac{\lambda}{2}\right)^2 - \frac{\lambda^2}{4} - 2\lambda + \lambda + (y-2)^2 = \left(x-1-\frac{1}{2}(\lambda+2)\right)^2 - \frac{1}{4}(\lambda+2)^2 + 1 + (y-2)^2$$

For $(x, y, \lambda) = (1, 2, -2)$ all terms in brackets vanish and this is 1, as it should be. All small excursions of x (or of y, or of λ) away from (1, 2, -2) yield quadratic changes: we have an stationary point. As we know, because the 3D gradient vanishes.

This stationary point is a global minimum on the xy-plane (with $\lambda = -2$ fixed). It is a saddle point in the 3D xy λ -space. Note: the method provides no guarantees about the nature of the stationary points.

Now in more, i.e., *n* dimensions (see Groenenboom notes). We have a function

$$f(\mathbf{x}) = f(x_1, x_2, \ldots, x_n)$$

We have *m* constraints:

$$g_i(\mathbf{x}) = g_i(x_1, x_2, \dots, x_n) = 0, \ i = 1, 2, \dots, m$$

This is a (n - m)-dimensional surface M.

We are looking for \mathbf{x}_e such that f is stationary in \mathbf{x}_e and on all $g_i(\mathbf{x}_e) = 0$.

Any allowed (by the constraints) excursion s away from \mathbf{x}_e is perpendicular to all gradients ∇g_i (calculated in \mathbf{x}_e):

$$\mathbf{s} \perp \nabla g_1$$
 and $\mathbf{s} \perp \nabla g_2$ and $\mathbf{s} \perp \nabla g_3, \dots, \mathbf{s} \perp \nabla g_m$. (1)

If excursions **s** are perpendicular to ∇f then there is *no* first-order change of *f*.

We can achieve this by demanding that the gradient ∇f be a linear combination of the gradients ∇g_i , i.e.

$$\nabla f = \sum_i \lambda_i \nabla g_i.$$

Then, by construction, any vector **s** (simultaneously on all constraints) is perpendicular to ∇f and gives a vanishing first-order change in f: \mathbf{x}_e is a stationary point.

We define (called $\tilde{f}(\vec{r}, \vec{\lambda})$ in Groenenboom):

$$\mathcal{L}(\mathbf{x}, oldsymbol{\lambda}) \equiv f(\mathbf{x}) - \sum_{i=1}^m \lambda_i g_i(\mathbf{x})$$

and we demand that the gradient in (n + m)-dimensional space vanishes:

$$\frac{\partial \mathcal{L}}{\partial x_j} = \frac{\partial f}{\partial x_j} - \sum_{i=1}^m \lambda_i \frac{\partial g_i}{\partial x_j} = 0 , \quad j = 1, 2, \dots, n$$
$$\frac{\partial \mathcal{L}}{\partial \lambda_i} = g_i = 0 , \quad i = 1, 2, \dots, m$$

Note:

- We are allowed to replace λ_i with $-\lambda_i$.
- We did not worry about complex x, ...
- ▶ Demanding $g_i(\mathbf{x}) = 0$ is general. Suppose $g_i(\mathbf{x}) = c \neq 0$, then $\tilde{g}_i(\mathbf{x}) = g_i(\mathbf{x}) c = 0$.

Why do we need Lagrange multipliers?

Because we have to minimize our energy expression:

$$\epsilon = 2\sum_{k}^{N/2} \langle \phi_k | \hat{h} | \phi_k \rangle + \sum_{k}^{N/2} \sum_{l}^{N/2} 2J_{kl} - K_{kl} = 2\sum_{k}^{N/2} h_{kk} + \sum_{k}^{N/2} \sum_{l}^{N/2} 2J_{kl} - K_{kl}$$

with the constraints that the orbitals are orthonormal, i.e.

$$\langle \phi_k | \phi_l \rangle = \delta_{lk}$$

These conditions make that the wave function, i.e., the Slater determinant, is normalized and satisfies Pauli.

In order to fulfill our constraints, we introduce Lagrange multipliers Λ_{ij} , one for each constraint, and add terms for our constraints to ϵ to make a new function ϵ' that we will minimize:

$$\epsilon' = 2\sum_{k}^{N/2} h_{kk} + \sum_{k}^{N/2} \sum_{l}^{N/2} (2J_{kl} - K_{kl}) - \sum_{k}^{N/2} \sum_{l}^{N/2} \Lambda_{lk} \left(\langle \phi_k | \phi_l \rangle - \delta_{lk} \right)$$

Towards Hartree-Fock: the basis

Help slide:

So we have to minimize ϵ' by changing the orbitals $|\phi_k\rangle$ while keeping the constraint $\langle \phi_k | \phi_l \rangle = \delta_{lk}$. This is quite abstract, and not suitable for a computer.

In the next slide we will expand the obitals in basis functions. The basis functions are **fixed**. The expansion coefficients $C_{\lambda k}$ are the **variables**. We need to find those values of $C_{\lambda k}$ that give the minimum of ϵ' .

In principle the number of basis functions M can be infinite. The number of orbitals is uniquely defined. In our closed shell system it is N/2, i.e. the number of electrons/2.

Towards Hartree-Fock: the basis

Expand the spatial orbitals in a basis:

$$\phi_k(\mathbf{r}) = \sum_{\lambda}^M C_{\lambda k} \, \boldsymbol{\chi}_{\lambda}(\mathbf{r}) = \sum_{\lambda}^M \, \boldsymbol{\chi}_{\lambda}(\mathbf{r}) C_{\lambda k} \quad \text{i.e.} \quad |\phi_k\rangle = \sum_{\lambda}^M |\boldsymbol{\chi}_{\lambda}\rangle C_{\lambda k}$$

Typically the ϕ_k are the orbitals of the whole system, i.e. the molecular orbitals (MOs). In chemistry the χ_{λ} are often chosen to be centered on the individual atoms: atomic orbitals (AOs). The acronym for this approximation: MO-LCAO (MO-Linear Combination of AOs).

In practice the basis is never complete, but should "approximate completeness sufficiently".

The atomic orbitals are, in general, not orthonormal, because

$$S_{\lambda\mu} = \langle \chi_{\lambda} | \chi_{\mu} \rangle \neq 0$$

if χ_{λ} and χ_{μ} are on different atoms. **S** is the overlap matrix.

In rare instances linear dependencies arise (a basis function $|\chi_{\lambda}\rangle$ is a linear combination of several others). The program will crash. That should be avoided.

Towards Hartree-Fock: the basis

We should minimize: $\epsilon' = 2 \sum_{k}^{N/2} h_{kk} + \sum_{k}^{N/2} \sum_{l}^{N/2} (2J_{kl} - K_{kl}) - \sum_{k}^{N/2} \sum_{l}^{N/2} \Lambda_{lk} (\langle \phi_k | \phi_l \rangle - \delta_{lk})$

We insert the expansion in the basis (using $|\phi_k\rangle = \sum_{\lambda}^{m} |\chi_{\lambda}\rangle C_{\lambda k}$):

$$\epsilon' = 2\sum_{k}^{N/2}\sum_{\lambda}^{M}\sum_{\mu}^{M}C_{\lambda k}^{*}C_{\mu k}\langle\chi_{\lambda}|\hat{h}|\chi_{\mu}
angle$$

$$+\sum_k^{N/2}\sum_l^{N/2}\sum_{\lambda}^M\sum_{\mu}\sum_{
u}^M\sum_{
u}^M\sum_{ au}^M\sum_{ au}^M C^*_{\lambda k}C^*_{\mu l}C_{
u k}C_{ au l}\langle\chi_\lambda\chi_\mu|(2-\hat{P}_{12})\hat{g}(1,2)|\chi_
u\chi_ au
angle \ -\sum_{
u}^{N/2}\sum_{\lambda}^{N/2}\Lambda_{lk}\left(\sum_{
u}^M\sum_{\lambda}^MC^*_{\lambda k}C_{\mu l}\langle\chi_\lambda|\chi_\mu
angle-\delta_{lk}
ight)$$

 $\begin{array}{c} \sum_{k} \sum_{l} \sum_{\mu} \sum_{\mu$ N/2 N/2

 $=\sum_{l=1}^{r}\sum_{l=1}^{r}2\langle\phi_{k}(1)\phi_{l}(2)|\hat{g}(1,2)|\phi_{k}(1)\phi_{l}(2)\rangle-\langle\phi_{k}(1)\phi_{l}(2)|\hat{P}_{12}\hat{g}(1,2)|\phi_{k}(1)\phi_{l}(2)\rangle$

Towards Hartree-Fock: "simplification"

Let's try to safe ink:

$$\begin{aligned} \epsilon' &= 2\sum_{k}^{N/2} \sum_{\lambda}^{M} \sum_{\mu}^{M} C_{\lambda k}^{*} C_{\mu k} \langle \chi_{\lambda} | \hat{h} | \chi_{\mu} \rangle \\ &+ \sum_{k}^{N/2} \sum_{l}^{N/2} \sum_{\lambda}^{M} \sum_{\mu}^{M} \sum_{\nu}^{M} \sum_{\tau}^{M} C_{\lambda k}^{*} C_{\mu l}^{*} C_{\nu k} C_{\tau l} \langle \chi_{\lambda} \chi_{\mu} | (2 - \hat{P}_{12}) \hat{g}(1, 2) | \chi_{\nu} \chi_{\tau} \rangle \\ &- \sum_{k}^{N/2} \sum_{l}^{N/2} \Lambda_{l k} \left(\sum_{\lambda}^{M} \sum_{\mu}^{M} C_{\lambda k}^{*} C_{\mu l} \langle \chi_{\lambda} | \chi_{\mu} \rangle - \delta_{i j} \right) \\ &= \sum_{\lambda \mu} 2 P_{\mu \lambda} h_{\lambda \mu} + \sum_{\lambda \mu \nu \tau} P_{\nu \lambda} P_{\tau \mu} \langle \lambda \mu | (2 - \hat{P}_{12}) \hat{g} | \nu \tau \rangle - \sum_{k l}^{N} \Lambda_{l k} (\sum_{\lambda \mu} C_{\lambda k}^{*} C_{\mu l} S_{\lambda \mu} - \delta_{l k}) \end{aligned}$$

The density matrix **P** and (remember) overlap matrix **S** are:

$$P_{\mu\lambda} \equiv \sum_{k}^{N/2} C_{\mu k} C^*_{\lambda k} \quad , \quad S_{\lambda\mu} \equiv \langle \chi_{\lambda} | \chi_{\mu} \rangle$$

Note the indices for \mathbf{P} , the second has the c.c. These are not operators!

At the minimum all variations of ϵ' with respect to the real and complex parts of C_{om} must vanish:

$$\frac{\partial \epsilon'}{\partial \text{Re}[C_{\rho m}]} = 0 \text{ and } \frac{\partial \epsilon'}{\partial \text{Im}[C_{\rho m}]} = 0 \text{ where } C_{\rho m} = \text{Re}[C_{\rho m}] + i\text{Im}[C_{\rho m}]$$

This yields the same as requireing:

$$\frac{\partial \epsilon'}{\partial C_{\rho m}} = 0$$
 and $\frac{\partial \epsilon'}{\partial C^*_{\rho m}} = 0$ where $C_{\rho m}$ and $C^*_{\rho m}$ are considered *independent*

Note: we single out a specific element $C_{\rho m}$ or $C^*_{\rho m}$ in the sum:

$$\frac{\partial P_{\mu\lambda}}{\partial C_{\rho m}} = \frac{\partial}{\partial C_{\rho m}} \left(\sum_{k}^{N/2} C_{\mu k} C_{\lambda k}^{*} \right) = \sum_{k}^{N/2} \delta_{\mu \rho} \delta_{m k} C_{\lambda k}^{*} = \delta_{\mu \rho} C_{\lambda m}^{*}$$
$$\frac{\partial P_{\mu\lambda}}{\partial C_{\rho m}^{*}} = \frac{\partial}{\partial C_{\rho m}^{*}} \left(\sum_{k}^{N/2} C_{\mu k} C_{\lambda k}^{*} \right) = \sum_{k}^{N/2} \delta_{\lambda \rho} \delta_{m k} C_{\mu k} = \delta_{\lambda \rho} C_{\mu m}$$

$$\begin{aligned} \frac{\partial P_{\mu\lambda}}{\partial C_{\rho m}} &= \delta_{\mu\rho} C_{\lambda m}^{*} \text{ and } \frac{\partial P_{\mu\lambda}}{\partial C_{\rho m}^{*}} = \delta_{\lambda\rho} C_{\mu m} \\ \frac{\partial}{\partial C_{\rho m}^{*}} \left(\sum_{\lambda\mu} 2 P_{\mu\lambda} h_{\lambda\mu} \right) &= \sum_{\lambda\mu} 2 \delta_{\lambda\rho} C_{\mu m} h_{\lambda\mu} = \sum_{\mu} 2 C_{\mu m} h_{\rho\mu} = \sum_{\nu} 2 h_{\rho\nu} C_{\nu m} \\ \frac{\partial}{\partial C_{\rho m}^{*}} \left(\sum_{\lambda\mu\nu\tau} P_{\nu\lambda} P_{\tau\mu} \langle \lambda\mu | (2 - \hat{P}_{12}) \hat{g} | \nu\tau \rangle \right) \end{aligned}$$
$$= \sum_{\lambda\mu\nu\tau} \delta_{\lambda\rho} C_{\nu m} P_{\tau\mu} \langle \lambda\mu | (2 - \hat{P}_{12}) \hat{g} | \nu\tau \rangle + \sum_{\lambda\mu\nu\tau} P_{\nu\lambda} \delta_{\mu\rho} C_{\tau m} \langle \lambda\mu | (2 - \hat{P}_{12}) \hat{g} | \nu\tau \rangle \\ &= \sum_{\mu\nu\tau} P_{\tau\mu} \langle \rho\mu | (2 - \hat{P}_{12}) \hat{g} | \nu\tau \rangle C_{\nu m} + \sum_{\lambda\nu\tau} P_{\nu\lambda} \langle \lambda\rho | (2 - \hat{P}_{12}) \hat{g} | \nu\tau \rangle C_{\tau m} \\ &= \sum_{\mu\nu\tau} P_{\tau\mu} \langle \rho\mu | (2 - \hat{P}_{12}) \hat{g} | \nu\tau \rangle C_{\nu m} + \sum_{\mu\tau\nu} P_{\tau\mu} \langle \mu\rho | (2 - \hat{P}_{12}) \hat{g} | \tau\nu \rangle C_{\nu m} \\ &= 2 \sum_{\nu} \sum_{\mu\tau} P_{\tau\mu} \langle \rho\mu | (2 - \hat{P}_{12}) \hat{g} | \nu\tau \rangle C_{\nu m} \end{aligned}$$

$$\frac{\partial}{\partial C_{\rho m}^{*}} \left(-\sum_{kl} \Lambda_{lk} (\sum_{\lambda \mu} C_{\lambda k}^{*} C_{\mu l} S_{\lambda \mu} - \delta_{lk}) \right)$$
$$= -\sum_{l} \sum_{\mu} \Lambda_{lm} C_{\mu l} S_{\rho \mu} = -\sum_{l} \sum_{\mu} S_{\rho \mu} C_{\mu l} \Lambda_{lm}$$

Here we have defined the Fock matrix **F**. Further we have the overlap matrix **S**, the coefficient matrix **C** and the Lagrange multiplier matrix **A**. **F**, **S** and **A** are Hermitian:

$$\mathbf{F}^{\dagger}=\mathbf{F}$$
 i.e. $F_{\mu
ho}^{*}=F_{
ho\mu}$, $\mathbf{S}^{\dagger}=\mathbf{S}$ and $\mathbf{\Lambda}^{\dagger}=\mathbf{\Lambda}$

Help slide:

The elements of the Fock matrix:

$$\mathcal{F}_{oldsymbol{
ho}
u} = oldsymbol{h}_{
ho
u} + \sum_{\mu au} oldsymbol{P}_{ au\mu} \langle
ho\mu| (2 - \hat{P}_{12}) \hat{g} |
u au
angle$$

So its size is $M \times M$.

Note it depends on the coefficient matrix **C** via the elements of the density matrix $P_{\tau\mu}$.