Example exam CTC2, NWI-MOL176

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- Put your **name**, student number, and date **on every page**. Also number your pages.
- A printed version of the **lecture notes can be used** during the exam.
- Include your scratch paper.

Question 1: Spin-orbit coupling in the NH radical

The NH radical has a triplet ground state, i.e., the electronic spin quantum number is S = 1. In the lecture we did not discuss electronic spin, but in the questions you can use that the commutation relations for the electron spin operators, $\hat{\boldsymbol{S}} = (\hat{S}_x, \hat{S}_y, \hat{S}_z)$, are exactly the same as for the orbital angular momentum operators $\hat{\boldsymbol{l}} = (\hat{l}_x, \hat{l}_y, \hat{l}_z)$. The electron spin functions are denoted as $|S, M_S\rangle$, analogous to the $|lm\rangle$ angular momentum functions. The projection quantum numbers M_S and m correspond to the same "space-fixed" z-axis.

The Hamiltonian (\hat{H}) for NH, neglecting vibration, is given by

$$\hat{H} = \hat{H}_{\rm rot} + \hat{H}_{\rm SO},$$

with rotational Hamiltonian (with rotational constant B),

$$\hat{H}_{\rm rot} = B\hat{l}^2,$$

and spin-orbit coupling

$$\hat{H}_{\rm SO} = A\hat{l}\cdot\hat{S}$$

where A is the spin-orbit coupling constant. The solutions Ψ_i with energy E_i of the Schrödinger equation

$$\hat{H}\Psi_i = E_i\Psi_i$$

can be expanded in an "uncoupled basis" as

$$\Psi_i = \sum_{lmM_S} |lm\rangle |SM_S\rangle c_{lmSM_S;i},$$

where $c_{lmSM_S;i}$ are the expansion coefficients.

- **1a**. Give the matrix elements of rotational Hamiltonian in the uncoupled basis, and give the possible values of all quantum numbers.
- 1b. Derive the matrix elements of the spin-orbit Hamiltonian in the uncoupled basis.
- **1c.** Which of quantum numbers l, m, S, and M_S are good quantum numbers? Briefly explain your answer.

The Hamiltonian \hat{H} is invariant under rotation. This makes it possible to define a new basis with additional good quantum numbers.

1d. Define such a basis and give all good quantum numbers for this basis. Also give the possible values for the quantum numbers in this basis.

If the molecule is put in a magnetic field along the x-axis, it is more convenient to express the wave function in a basis where the projection quantum numbers refer to the projection of the angular momentum on the y-axis, rather than the z-axis. Thus, we define new angular momentum states,

$$|lmSM_S\rangle^{(x)} = \hat{R} \ (|lm\rangle|SM_S\rangle).$$

1e. Give the expression for the rotation operator \hat{R} that rotates the (rotational) wave function around the *y*-axis, such that in the states $|lmSM_S\rangle^{(x)}$ the projection of the angular momentum on the *x*-axis is $\hbar m$ and give the matrix elements of \hat{R} in the space-fixed basis with functions $|lmSM_S\rangle$.

Question 2: Parity in body-fixed coordinates

Consider an atom-diatom (A–BC) system in Jacobi coordinates (\mathbf{r}, \mathbf{R}) , where

$$oldsymbol{r} \equiv oldsymbol{r}_{
m B} - oldsymbol{r}_{
m C} \ oldsymbol{R} \equiv oldsymbol{r}_{
m A} - oldsymbol{r}_{
m BC}$$

and \mathbf{r}_{BC} is the center-of-mass of molecule BC. The inversion operator \hat{i} acts on the atomic coordinates \mathbf{r}_{X} with X = A, B, or C as

$$\hat{i}m{r}_{\mathrm{X}} = -m{r}_{\mathrm{X}}$$

2a. Derive the action of \hat{i} on the Jacobi coordinates r and R (give all steps of the derivation).

We define a new "body-fixed" coordinate frame, with the z-axis b_z along the Jacobi vector,

$$\boldsymbol{b}_z \equiv \frac{1}{R} \boldsymbol{R}, \quad (\text{with } R = |\boldsymbol{R}|).$$

The x-axis of the body-fixed frame is define by the cross product

$$\boldsymbol{b}_x \equiv |N| \, \boldsymbol{e}_z \times \boldsymbol{b}_z,$$

where N normalizes the vector \boldsymbol{b}_x , and \boldsymbol{e}_z is the space-fixed z-axis

$$\boldsymbol{e}_z = \begin{pmatrix} 0\\ 0\\ 1 \end{pmatrix}$$
 .

Finally, the body-fixed y axis is defined by

$$\boldsymbol{b}_y = \boldsymbol{b}_z \times \boldsymbol{b}_x,$$

such that the matrix

$$\boldsymbol{B} = [\boldsymbol{b}_x \boldsymbol{b}_y \boldsymbol{b}_z]$$

defines a right-handed orthonormal frame.

2b. Derive the action of the inversion operator \hat{i} onto the body-fixed frame **B**.

The diatomic coordinate r can be expressed in the body-fixed frame by

$$\boldsymbol{r} = \boldsymbol{B}\boldsymbol{r}^{(\mathrm{BF})}.$$

- 2c. Using the action of the inversion operator \hat{i} on the vectors \boldsymbol{r} and the columns of the matrix \boldsymbol{B} , derive the action of \hat{i} on the three components of the body-fixed vector $\boldsymbol{r}^{(B)}$
- 2d. Give the matrix-representation of the inversion operator \hat{i} in the body-fixed frame B, and show that the matrix corresponds to a reflection and give the reflection plane.

Question 3: The HF molecule in an electric field

The interaction of HF, which has an electric dipole moment of $\mu = 0.7 \ ea_0$, with an electric field of strength *E* along the space-fixed *z*-axis is given by the Stark Hamiltonian

$$\hat{H}_{\text{Stark}} = -E\mu\cos(\theta),$$

where θ is the angle between the molecular axis and the space-fixed z-axis. Here we consider the molecule HF with rotational states $|lm\rangle$.

3a. Derive an expression for the matrix elements of the Stark Hamiltonian in the rotational basis $|lm\rangle$,

$$H_{l'm';lm} = \langle l'm' | \hat{H}_{\text{Stark}} | lm \rangle.$$

3b. Give the selection rules for these matrix elements, i.e., for which combinations of the quantum numbers l', m', l, and m are the matrix elements equal to zero?

The Hamiltonian that describes the rotation of HF (with dipole moment μ) in an electric field (of strength E), is given by

$$\hat{H} = B\hat{l}^2 + \hat{H}_{\text{Stark}}.$$

3c. Write down the time-dependent Schrödinger equation for this system and give the matrix elements of \hat{H} in the rotational basis.

Consider a molecule that is in l = 0 at time t = 0. For simplicity, we only consider the coupling by the electric field of the initial state to the rotational state $|l, m\rangle = |1, 0\rangle$.

3d. Solve the time-dependent Schrödinger equation, and give the probability P(t) of observing the molecule in the l = 1 state at times t > 0.